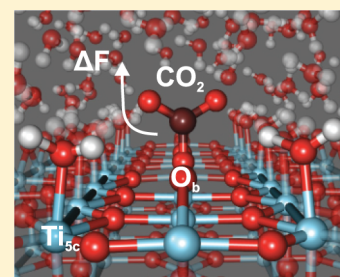


# CO<sub>2</sub> Adsorption and Reactivity on Rutile TiO<sub>2</sub>(110) in Water: An *Ab Initio* Molecular Dynamics Study

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## Supporting Information

**ABSTRACT:** Atomic-scale understanding of CO<sub>2</sub> adsorption and reactivity on TiO<sub>2</sub> is important for the development of new catalysts for CO<sub>2</sub> conversion with improved efficiency and selectivity. Here, we employ Car–Parrinello molecular dynamics combined with metadynamics simulations to explore the interaction dynamics of CO<sub>2</sub> and rutile TiO<sub>2</sub>(110) surface explicitly treating water solution at 300 K. We focus on understanding the competitive adsorption of CO<sub>2</sub> and H<sub>2</sub>O, as well as the kinetics of CO and bicarbonate (HCO<sub>3</sub><sup>−</sup>) formation. Our results show that adsorption configurations and possible reaction pathways are greatly affected by proper description of the water environment. We find that in aqueous solution, CO<sub>2</sub> preferentially adsorbs at the bridging oxygen atom O<sub>b</sub>, while Ti<sub>5c</sub> sites are saturated by H<sub>2</sub>O molecules that are difficult to displace. Our calculations predict that further conversion reactions include spontaneous protonation of adsorbed CO<sub>2</sub> and detachment of OH<sup>−</sup> to form a CO molecule that is significantly facilitated in the presence of a surface Ti<sup>3+</sup> polaron. In addition, the mechanisms of HCO<sub>3</sub><sup>−</sup> formation in bulk water and near TiO<sub>2</sub>(110) surface are discussed. These results provide atomistic details on the mechanism and kinetics of CO<sub>2</sub> interaction with TiO<sub>2</sub>(110) in a water environment.



## INTRODUCTION

The problem of anthropogenic CO<sub>2</sub> emissions has been recently attracting a lot of attention due to a dramatic rise of CO<sub>2</sub> in the atmosphere in the past few decades.<sup>1–3</sup> One efficient way to counterbalance CO<sub>2</sub> emissions is to photocatalytically convert CO<sub>2</sub> to valuable chemicals such as methanol (CH<sub>3</sub>OH), methane (CH<sub>4</sub>), formic acid (HCOOH), and formaldehyde (H<sub>2</sub>CO) using renewable solar energy.<sup>2–6</sup> Among various photocatalysts, anatase and rutile phases of titanium dioxide (TiO<sub>2</sub>) have been extensively examined both experimentally and theoretically. Despite a great deal of prior research,<sup>5,7–15</sup> both the efficiency and selectivity of CO<sub>2</sub> conversion reactions over TiO<sub>2</sub> catalysts are still poor, and the reaction mechanisms are not yet well understood at the atomic level.

Many catalytic CO<sub>2</sub> conversion reactions are experimentally carried out in aqueous solution, while previous theoretical investigations have been limited to consideration of CO<sub>2</sub> reaction mechanisms in the gas phase<sup>7–10</sup> and primarily for the anatase polymorph of TiO<sub>2</sub>.<sup>15,16–18</sup> It is well established, however, that the presence of a solvent can play a critical role in determining both the rates and pathways of various reactions including CO<sub>2</sub> conversion. For example, it is known that methanol is the major product in the gas-phase reduction of CO<sub>2</sub> on Cu catalysts, whereas methane is the predominant species in the aqueous electrocatalytic reduction process.<sup>19,20</sup>

Some recent first-principles computational studies<sup>7,9</sup> have indeed demonstrated that the presence of even one H<sub>2</sub>O molecule coadsorbing with CO<sub>2</sub> on different TiO<sub>2</sub> surfaces can change the binding energy of CO<sub>2</sub> by as much as 0.25 eV

relative to the case with no water.<sup>7</sup> Moreover, not only can the stability of CO<sub>2</sub> molecule on the TiO<sub>2</sub> surfaces be modified by water environment, but also the stability and reactivity of reaction intermediates. Although some of the previous theoretical works have explored the effect of H<sub>2</sub>O molecule coadsorption on the reaction mechanisms of surface-mediated CO<sub>2</sub> conversion, most of the studies have been limited to introducing the aqueous environment through the implicit solvent model. This model, however, cannot account for the formation and dynamics of hydrogen bonding network that greatly affect CO<sub>2</sub> adsorption and surface-mediated reaction mechanisms.<sup>19</sup>

In recent years, computational techniques such as metadynamics that allow one to avoid the limitations of traditional *ab initio* molecular dynamics (AIMD) to sample the free energy landscape have been gaining much attention.<sup>9,16,19,21</sup> Recently, AIMD-based metadynamics simulations were successfully employed to unveil the mechanisms and kinetics of various interfacial reactions including electrochemical reduction of CO<sub>2</sub> and CO in aqueous solutions at room temperature over model catalyst surfaces such as the Cu(100) facet.<sup>19,22,23</sup>

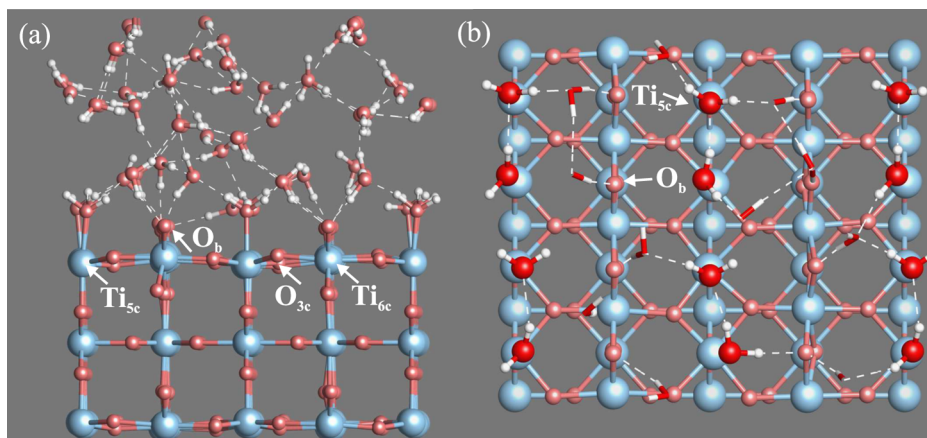
## COMPUTATIONAL METHODOLOGY

In this study, we carry out density-functional-theory (DFT) Car–Parrinello molecular dynamics (CPMD) simulations to allow explicit treatment of water and temperature effects on

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**Figure 1.** Side (a) and top (b) views of the interfacial water structure at the rutile  $\text{TiO}_2(110)$  surface from CPMD simulations. Color code: Ti, blue; O, red; H, white. The water molecules of the first layer are represented by a ball and stick model, the dashed lines highlight the hydrogen bonds network.

$\text{CO}_2$  adsorption and reactivity at the rutile  $\text{TiO}_2(110)$  surface. To obtain free-energy reaction barriers, we utilize a CPMD-based metadynamics technique that allows accelerated sampling of reaction pathways characterized by a set of collective variables (CVs), as implemented in the NWChem software package.<sup>24</sup> Previously, metadynamics simulations have been successfully applied to examine the mechanisms and rates of a variety of rare-event chemical reactions,<sup>9,16,19,21</sup> and the theoretical background of this approach was described in detail elsewhere.<sup>25,26</sup> Full details of our computational protocol can be found in the [Supporting Information](#), while here we just briefly describe the key features of the computational scheme.

The rutile  $\text{TiO}_2(110)$  surface was modeled as a periodic slab comprised of the three  $\text{TiO}_2$  layers (the central layer was fixed) with a  $4 \times 2$  surface supercell and a vacuum gap of 10 Å, thus resulting in a  $13.38 \times 11.91 \times 19.98 \text{ \AA}^3$  simulation cell. To provide a water density of about  $1 \text{ g/cm}^3$ , the vacuum gap was filled with one  $\text{CO}_2$  and 55  $\text{H}_2\text{O}$  molecules (see [Figure 1](#)). All calculations were performed at the  $\Gamma$  point using the Perdew–Burke–Ernzerhof (PBE) exchange–correlation functional,<sup>27</sup> Troullier–Martins<sup>28</sup> pseudopotentials for Ti, and Hamann-type<sup>29</sup> pseudopotentials for H, C, and O atoms. The kinetic energy cutoff of 60 Ry was employed, while a larger energy cutoff of 100 Ry was found to yield very similar  $\text{CO}_2$  desorption energetics (0.65 eV for the adsorption position  $\text{O}_b\text{-CO}_2$ , as compared with 0.67 eV at 60 Ry). The PBE functional was corrected for long-range dispersion interactions using the Grimme approach (DFT-D4).<sup>30</sup> The temperature of the simulation cell was maintained constant at 300 K employing a Nose–Hoover thermostat.<sup>31,32</sup> To achieve thermal stability, 6 ps quantum mechanics/molecular mechanics (QM/MM) simulations were first carried out to pre-equilibrate the water/ $\text{TiO}_2(110)$  interface followed by 6 ps of CPMD equilibration. Then each adsorption configuration was additionally equilibrated for  $\sim 5$  ps to provide energy conservation within  $10^{-3}$  eV/ps. For each adsorbed state, we used several slightly different starting configurations. The equilibration procedure used in this study is similar to some previous AIMD studies of water/metal oxide interfaces including  $\text{TiO}_2$ .<sup>16,19,21,33</sup> The CPMD production trajectories were 6 ps long, while metadynamics simulations were run for 4–12 ps depending on the system. For the reactions involving a surface  $\text{Ti}^{3+}$  polaron, we employ the GGA+ $U$  formalism<sup>34</sup> with  $U_{\text{eff}} = 5$

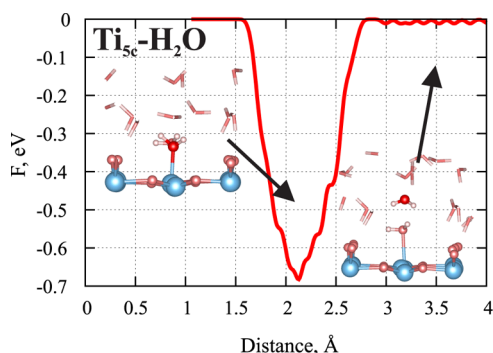
eV based on previous theoretical studies of polarons in  $\text{TiO}_2$ .<sup>35,36</sup>

## RESULTS AND DISCUSSION

**$\text{TiO}_2(110)$ /Water Interface.** We start by considering the structure of interfacial water and energetics of water desorption from the hydroxylated  $\text{TiO}_2(110)$  surface, which is important due to the competition between  $\text{H}_2\text{O}$  and  $\text{CO}_2$  molecules in the aqueous phase for surface adsorption. The molecular structure of water– $\text{TiO}_2(110)$  interface was previously examined in a number of experimental<sup>37–40</sup> and first-principles computational<sup>37,40–44</sup> investigations, and our CPMD results reported here are in full agreement with a recent detailed study by *in situ* scanning probe microscopy.<sup>40</sup> The rutile  $\text{TiO}_2(110)$  surface layer contains the 5- and 6-fold coordinated Ti atoms ( $\text{Ti}_{5c}$  and  $\text{Ti}_{6c}$ , respectively), 3-fold coordinated O atoms ( $\text{O}_{3c}$ ) and 2-fold coordinated bridging O atoms ( $\text{O}_b$ ), as depicted in [Figure 1](#). As an initial state for all calculations the uppermost surface of the slab was saturated by  $\text{H}_2\text{O}$  molecules attached to the undercoordinated surface  $\text{Ti}_{5c}$  cations, while the rest of the vacuum gap was filled randomly with  $\text{H}_2\text{O}$  molecules before starting QM/MM followed by CPMD simulations.

We did not observe any water dissociation over the  $\text{TiO}_2$  surface during equilibration at 300 K, in accordance with previous first-principles modeling studies.<sup>42,43</sup> For the equilibrated state, we find that the first interfacial layer of water along the  $\text{Ti}_{5c}$  rows exhibits molecular ordering characterized by a water–dimer structural motif as shown in [Figure 1](#), in full agreement with a recent first-principles molecular dynamics investigation.<sup>40</sup> Such interfacial structure is different from a one-dimensional chain structure along the  $\text{Ti}_{5c}$  rows suggested by static DFT calculations<sup>44</sup> and is due to the interaction between  $\text{H}_2\text{O}$  molecules of the first and second water layers through the hydrogen bonding network. Specifically, we identify a stable network of hydrogen bonds between  $\text{O}_b$  atoms, adsorbed  $\text{H}_2\text{O}$  at the  $\text{Ti}_{5c}$  sites, and the  $\text{H}_2\text{O}$  molecules in the second hydration layer. The third and farther water layers do not interact directly with the  $\text{TiO}_2$  surface and can be considered as the bulk liquid phase.

To analyze the energetics of the  $\text{H}_2\text{O}$  desorption process, we only consider desorption of  $\text{H}_2\text{O}$  from the  $\text{Ti}_{5c}$  site, since  $\text{H}_2\text{O}$  molecules interacting with other surface sites through the hydrogen bonds are expected to be labile at 300 K. [Figure 2](#)

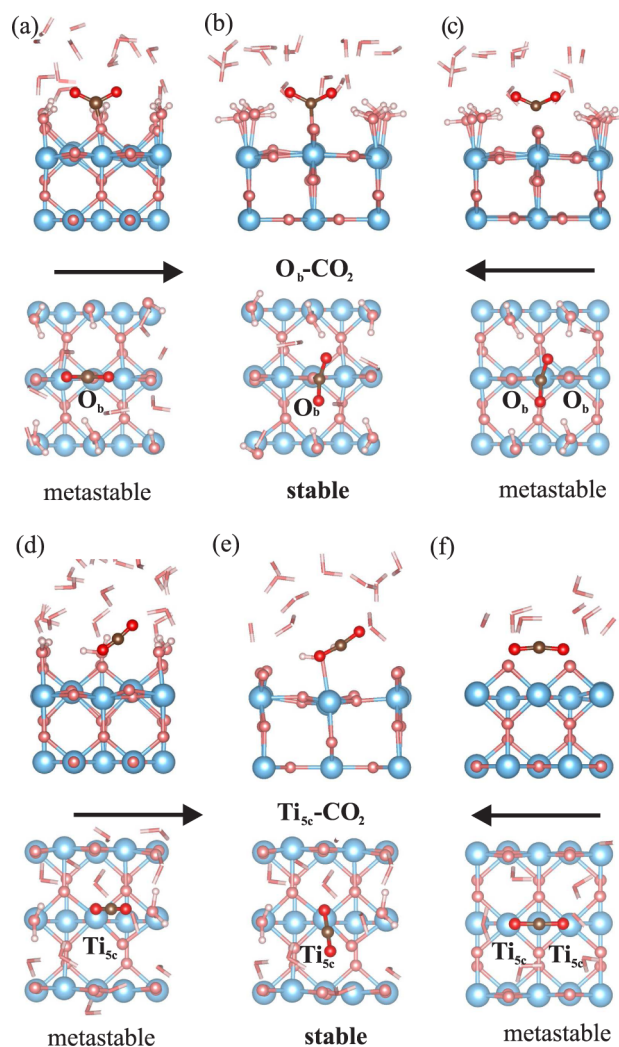


**Figure 2.** Free-energy profile of H<sub>2</sub>O desorption reaction from the Ti<sub>5c</sub> surface site as computed using CPMD metadynamics simulations with the distance (in Å) between adsorbed H<sub>2</sub>O and TiO<sub>2</sub>(110) surface used as a collective variable. The initial and final states are shown on the left and right, respectively.

shows the free-energy profile of the desorption reaction based on CPMD metadynamics simulations as a function of the bond distance between Ti<sub>5c</sub> and O<sub>w</sub> of the H<sub>2</sub>O molecule leaving the surface. It is seen that when the bond distance becomes larger than 3 Å, the desorption reaction can be considered complete. The activation barrier for H<sub>2</sub>O desorption is estimated at 0.69 eV, thus being in the experimental range of 0.6–1.05 eV,<sup>10,45</sup> but it is smaller than the values of 0.77 eV<sup>42</sup> and 0.99 eV<sup>9</sup> estimated previously from DFT with no explicit treatment of water at 0 K. Our smaller value can be explained by additional stabilization of the desorbed H<sub>2</sub>O molecule due to the hydrogen bonding network of the bulk water.

**CO<sub>2</sub> Adsorption and Reduction.** We next examine different CO<sub>2</sub> adsorption configurations at the TiO<sub>2</sub>(110) surface in the presence of water at 300 K. Out of six adsorption states observed in previous static DFT calculations at 0 K,<sup>7,9</sup> only two are found to be stable after a few picoseconds of CPMD simulations, as shown in Figure 3. The first stable adsorption state is O<sub>b</sub>-CO<sub>2</sub> (Figure 3b), where CO<sub>2</sub> is adsorbed at the bridging oxygen atom O<sub>b</sub>. The average O<sub>b</sub>-C bond distance during 5 ps trajectory is 1.365 Å, while the O-C-O angle is about 124.7°. This is very close to that of the CO<sub>2</sub> molecule in the gas phase (138°)<sup>46</sup> and much smaller than that found for the adsorbed CO<sub>2</sub> in DFT-PBE calculations with no water at 0 K (178°).<sup>7,9</sup> This suggests that some activation has occurred owing to the electron transfer from the TiO<sub>2</sub> surface to the adsorbed CO<sub>2</sub> species. A number of other configurations such as CO<sub>2</sub> at the O<sub>b</sub> site rotated by 90° around the z axis relative to the O<sub>b</sub>-CO<sub>2</sub> configuration (Figure 3a) and CO<sub>2</sub> located between two adjacent bridging O<sub>b</sub> atoms (Figure 3c) switched to the O<sub>b</sub>-CO<sub>2</sub> configuration after a few picoseconds.

It is believed that electron transfer from the TiO<sub>2</sub> surface to adsorbed CO<sub>2</sub> activates the molecule toward further reactions via the formation of a bent CO<sub>2</sub> species.<sup>7,47</sup> It was estimated in static DFT calculations that a relatively low energy barrier of 0.27 eV is required to activate a planar physisorbed CO<sub>2</sub> molecule to a bent adsorbed configuration.<sup>9</sup> In our calculations, we observe that the physisorbed CO<sub>2</sub> molecule is weakly bound to the TiO<sub>2</sub>(110) surface and may spontaneously leave the surface during CPMD. We also see (Figure 1b) that the second layer of water bounded to O<sub>b</sub> may also hinder CO<sub>2</sub> physisorption. Thus, despite the predicted low barrier, CO<sub>2</sub> adsorption and further activation appears to be the rate-limiting step of CO<sub>2</sub> reduction.<sup>6,48</sup> It was also hypothesized that



**Figure 3.** Side and top views of the main adsorption configurations of CO<sub>2</sub> on the rutile TiO<sub>2</sub>(110) surface. (a–c) CO<sub>2</sub> adsorbed on a bridging oxygen (O<sub>b</sub>); (d–f) CO<sub>2</sub> adsorbed on the five-coordinated Ti atom (Ti<sub>5c</sub>).

photogeneration of an electron–hole pair might contribute to this activation; however, the role of this process in CO<sub>2</sub> activation is still unclear, and a DFT study of excited stoichiometric anatase TiO<sub>2</sub> surface has shown no tendency for electron transfer from the surface to adsorbed CO<sub>2</sub>.<sup>49</sup> Thus, our results suggest that CO<sub>2</sub> in aqueous solution is partially activated in the O<sub>b</sub>-CO<sub>2</sub> configuration; however, the presence of an electron–hole pair or intrinsic surface defects can additionally contribute to CO<sub>2</sub> activation.

The second stable adsorption state is Ti<sub>5c</sub>-CO<sub>2</sub>, which is a monodentate adsorbed CO<sub>2</sub> molecule bound to the Ti<sub>5c</sub> site via an oxygen of CO<sub>2</sub> in a tilted configuration (Figure 3e). The average bond distance Ti–O is equal to 2.674 Å, and the O–C–O angle is 173.2°. A similar 90° rotated configuration (Figure 3d) is found to be metastable and switched to Ti<sub>5c</sub>-CO<sub>2</sub>. The bidentate configuration of CO<sub>2</sub> adsorbed along the Ti<sub>5c</sub> row (Figure 3f) is also found to be metastable and changed to the Ti<sub>5c</sub>-CO<sub>2</sub> adsorption state.

The adsorption energies of CO<sub>2</sub> on TiO<sub>2</sub> surfaces are known to be greatly affected by introducing a water environment. Specifically, the binding energy of CO<sub>2</sub> was determined to be larger by ~0.1–0.2 eV when considering coadsorption with a

single H<sub>2</sub>O<sup>7,9</sup> and by ~0.3 eV when applying an implicit solvent model.<sup>7</sup> A comparison of DFT calculated adsorption energies for CO<sub>2</sub> on TiO<sub>2</sub> in various models is provided in Table 1.

**Table 1. Adsorption Energies (in eV) of a Single CO<sub>2</sub> Molecule at the O<sub>b</sub> and Ti<sub>5c</sub> Surface Sites with None (*E*<sub>ads</sub>) and Two Coadsorbing H<sub>2</sub>O Molecules (*E*<sub>coads</sub>), As Well As Employing Implicit Solvent Model (*E*<sub>ISM</sub>) Based on Zero-Temperature DFT Calculations<sup>a</sup>**

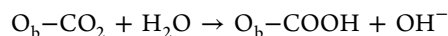
	<i>E</i> <sub>ads</sub>	<i>E</i> <sub>coads</sub>	<i>E</i> <sub>ISM</sub>
O <sub>b</sub> -CO <sub>2</sub>	0.22, <sup>7</sup> 0.29 <sup>9</sup>	0.63, 0.46, <sup>7</sup> 0.54 <sup>9</sup>	0.83 <sup>7</sup>
Ti <sub>5c</sub> -CO <sub>2</sub>	0.33, <sup>7</sup> 0.46 <sup>9</sup>	0.41, 0.42, <sup>7</sup> 0.49 <sup>9</sup>	0.75-0.81 <sup>7</sup>

<sup>a</sup>Adsorption energy of a single H<sub>2</sub>O molecule at Ti<sub>5c</sub> estimated as 1.01 eV is provided for comparison.

Overall, the results suggest that CO<sub>2</sub> is thermodynamically equally stable at the Ti<sub>5c</sub> and O<sub>b</sub> surface sites, whereas H<sub>2</sub>O is stronger adsorbed to the rutile surface (at the Ti<sub>5c</sub> site) than CO<sub>2</sub> (at any site). Thus, CO<sub>2</sub> molecules are expected to preferentially bind to the O<sub>b</sub> sites based on zero-temperature DFT results. To better understand the competitive adsorption/desorption between H<sub>2</sub>O and CO<sub>2</sub> at the rutile surface in water at 300 K, we perform additional free-energy calculations to estimate CO<sub>2</sub> desorption barriers and compare it with the H<sub>2</sub>O case.

Figure 4 shows the simulated free-energy profiles of CO<sub>2</sub> desorption from Ti<sub>5c</sub> and O<sub>b</sub> surface sites, while Table 2 summarizes the estimated desorption barriers for all species considered in this study. It is apparent that not only H<sub>2</sub>O has a higher adsorption energy than CO<sub>2</sub> (Table 1), but also it is much more difficult to desorb it from the Ti<sub>5c</sub> site (Table 2). On the other hand, it is considerably more energetically difficult to displace adsorbed CO<sub>2</sub> from the O<sub>b</sub> site than from Ti<sub>5c</sub> site. Therefore, our results suggest that in aqueous environment the Ti<sub>5c</sub> sites should be filled with H<sub>2</sub>O molecules, while CO<sub>2</sub> preferentially adsorbs on the O<sub>b</sub> sites.

For the most stable O<sub>b</sub>-CO<sub>2</sub> configuration we observe spontaneous protonation of the adsorbed CO<sub>2</sub> molecule during CPMD with a proton coming from the H<sub>2</sub>O molecule adsorbed at the neighboring Ti<sub>5c</sub> site according to the following reaction:



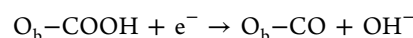
The driving force for such a spontaneous proton shuttling can be the activated character of \*COO (\*COO denotes adsorbed COO) in this configuration. Figure 5 illustrates how the adsorbed CO<sub>2</sub> molecule exchanges an H atom with the

**Table 2. Comparison of Free-Energy Barriers (in eV) of CO<sub>2</sub>, H<sub>2</sub>O, and HCO<sub>3</sub><sup>-</sup> Desorption from the O<sub>b</sub> and Ti<sub>5c</sub> Sites of the Rutile TiO<sub>2</sub>(110) Surface in Aqueous Solution from CPMD-Based Metadynamics Simulations**

O <sub>b</sub> -CO <sub>2</sub>	Ti <sub>5c</sub> -CO <sub>2</sub>	Ti <sub>5c</sub> -H <sub>2</sub> O	Ti <sub>5c</sub> -HCO <sub>3</sub> <sup>-</sup> (monodentate)	Ti <sub>5c</sub> -HCO <sub>3</sub> <sup>-</sup> (bidentate)
0.67	0.27	0.69	0.49	0.87

neighboring H<sub>2</sub>O molecules during 6.5 ps CPMD time frame. As seen from Figure 5 at any instance of time, only one (or none) of the two nearest-neighbor H atoms can be bound to the CO<sub>2</sub> molecule: when the first proton becomes bound, the second is repelled from CO<sub>2</sub>. When considering CO<sub>2</sub> desorption, however, CO<sub>2</sub> always leaves the surface with no protons attached. Thus, our simulations demonstrate that there is a dynamic equilibrium between adsorbed \*COO and \*COOH species at the O<sub>b</sub> surface site.

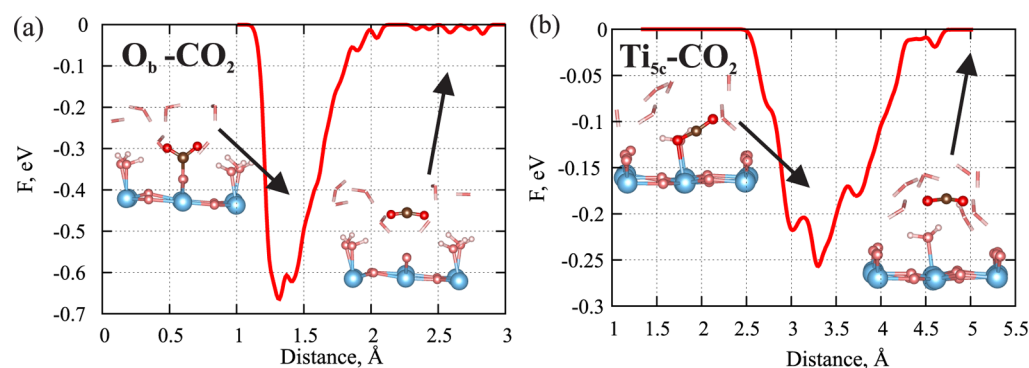
We next analyze the reduction of CO<sub>2</sub> to CO, which was experimentally shown to be the dominant reaction pathway on pure TiO<sub>2</sub>(110) surface.<sup>47</sup> According to a previously proposed mechanism,<sup>4,5,50</sup> we consider the following reaction of CO formation through dehydroxylation of \*COOH formed at the surface:



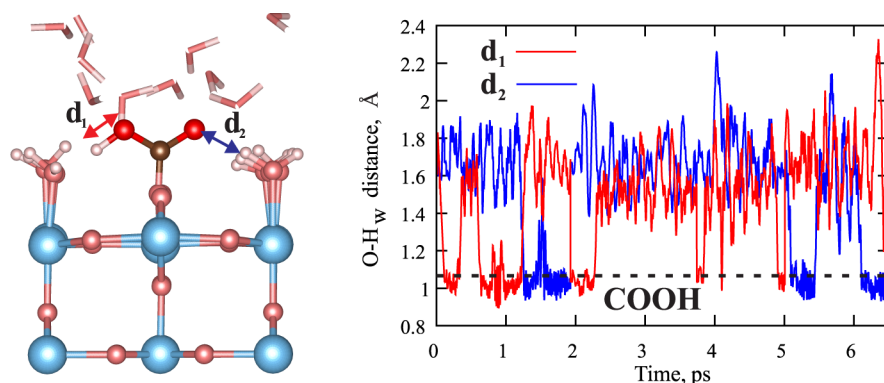
It was previously demonstrated in static DFT calculations that the surface Ti<sup>3+</sup> polaron is more stable at Ti<sub>5c</sub> than at Ti<sub>6c</sub> site.<sup>35,36</sup> Moreover, we find that Ti<sup>3+</sup> polaron is indeed stable at Ti<sub>5c</sub> during metadynamics simulations at 300 K, with an electron eventually being transferred to the \*COOH molecule followed by dehydroxylation reaction. On the contrary, an attempt to stabilize the polaron at Ti<sub>6c</sub> was not successful, and after a short period of time we observe spontaneous protonation of the surface O<sub>b</sub> atom by a nearby H<sub>2</sub>O molecule accompanied by electron delocalization from the Ti<sub>6c</sub> site.

Figure 6 compares the free-energy profiles of CO formation from \*COOH obtained with and without surface Ti<sup>3+</sup> polaron. Our calculations indicate that the presence of Ti<sup>3+</sup> polaron facilitates the dehydroxylation of COOH substantially by making the energy barrier of the reaction almost twice lower than without a polaron. This result is completely consistent with experimental observations of the absence of this reaction in the dark<sup>51</sup> and significant reaction acceleration on oxygen-deficient surfaces,<sup>48</sup> where the oxygen vacancies provide excess electrons to the rutile TiO<sub>2</sub>(110).

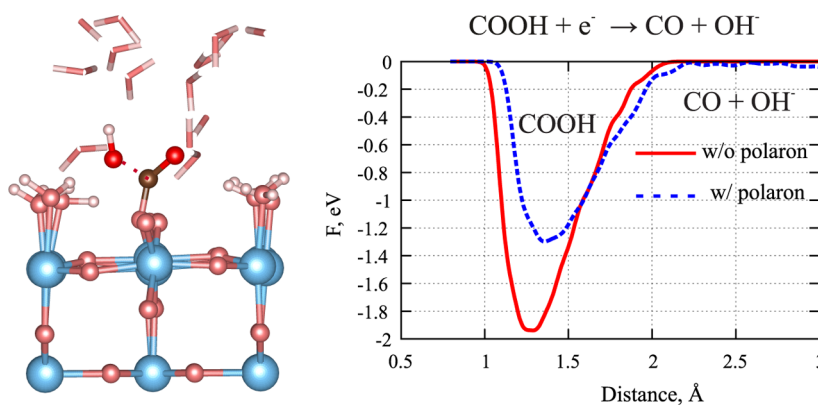
**Formation of HCO<sub>3</sub><sup>-</sup> Species and Their Stability.** The formation of bicarbonate HCO<sub>3</sub><sup>-</sup> species was detected in a



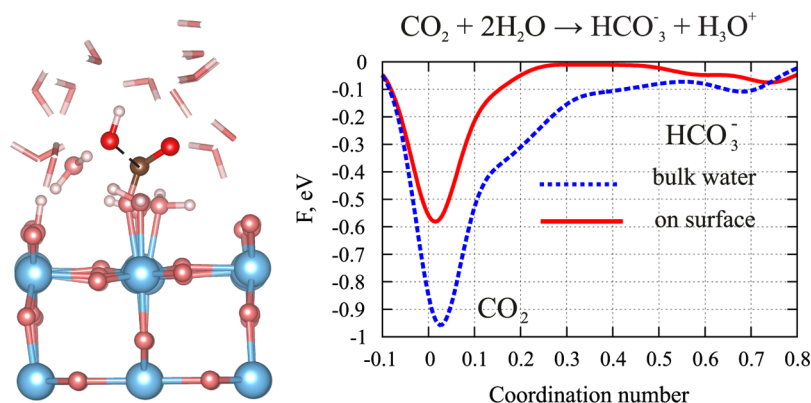
**Figure 4.** Free-energy profiles of CO<sub>2</sub> desorption reaction from (a) O<sub>b</sub> site and (b) Ti<sub>5c</sub> site of the TiO<sub>2</sub>(110) surface.



**Figure 5.** Illustration of CO<sub>2</sub> protonation/deprotonation reaction during 6.5 ps CPMD run with proton shuttling between adsorbed CO<sub>2</sub> and nearby H<sub>2</sub>O molecules. The O–H<sub>w</sub> distance less than 1.05 Å means that an H belongs to the CO<sub>2</sub> molecule.



**Figure 6.** Free-energy profiles of CO formation from the adsorbed COOH molecule with and without a surface Ti<sup>3+</sup> polaron as computed by CPMD metadynamics within the GGA+*U* approach.

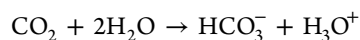


**Figure 7.** Free-energy profiles along the reaction pathway of HCO<sub>3</sub><sup>−</sup> formation near the TiO<sub>2</sub>(110) surface and in the bulk water. The energy profile is computed using the coordination number of a C atom with O atoms from water as a collective variable. The value of ~0.7 means that C is bound to OH, forming HCO<sub>3</sub><sup>−</sup>.

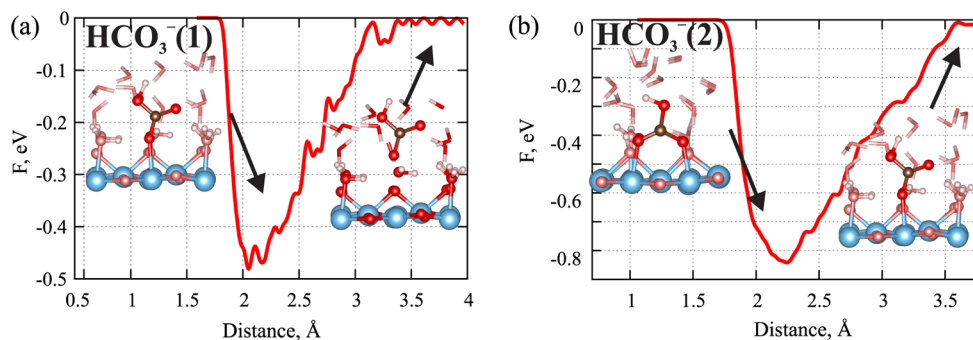
number of experiments,<sup>15,52</sup> while HCO<sub>3</sub><sup>−</sup> was not observed in some other experiments at the surface under ambient conditions.<sup>52</sup> It was also found that under supercritical conditions, almost all CO<sub>2</sub> in water solution can be transformed to HCO<sub>3</sub><sup>−</sup>.<sup>53</sup> HCO<sub>3</sub><sup>−</sup> formation is assumed to play an important role in the CO<sub>2</sub> conversion process in aqueous solution since the monolayer of bicarbonate can block the catalyst surface. Therefore, to gain insight into the mechanism of HCO<sub>3</sub><sup>−</sup> formation reaction both in the bulk water and at the TiO<sub>2</sub>(110) surface, we analyze HCO<sub>3</sub><sup>−</sup> stability on different

surface sites and the competition for these sites between HCO<sub>3</sub><sup>−</sup> and H<sub>2</sub>O adsorption/desorption.

The formation of HCO<sub>3</sub><sup>−</sup> species can proceed according to the following reaction:



In order to evaluate the catalytic effect of the TiO<sub>2</sub>(110) surface in bicarbonate formation, we compare the kinetics of this reaction in bulk water and at the rutile surface based on metadynamics simulations. In the bulk water case, the reaction was simulated in a periodic cubic box of length 13 Å filled with



**Figure 8.** Main adsorption configurations of  $\text{HCO}_3^-$  on the  $\text{TiO}_2(110)$  surface and the corresponding free-energy profiles for desorption reactions.

74  $\text{H}_2\text{O}$  molecules and one  $\text{CO}_2$  molecule using the same CPMD computational protocol as before. To drive the reaction, the number of water oxygen atoms bound to the carbon atom was used as a collective variable. We observe that in the proximity of the barrier crossing,  $\text{CO}_2$  becomes activated by bending toward a neighboring  $\text{H}_2\text{O}$  molecule, and after that the  $\text{H}_2\text{O}$  deprotonation occurs. The estimated free-energy barrier of the reaction in bulk water is about 0.85 eV (Figure 7), which is in excellent agreement with the experimental value (0.95 eV)<sup>54</sup> and previous CPMD calculations using a similar computational protocol (0.82 eV).<sup>21</sup>

For the most favorable adsorption configuration of  $\text{CO}_2$  on  $\text{TiO}_2$  ( $\text{O}_b\text{-CO}_2$ ) we could not drive  $\text{HCO}_3^-$  formation in metadynamics simulations, partially due to the unfavorable steric effect of this  $\text{CO}_2$  configuration. By contrast, the reaction at  $\text{Ti}_{5c}$  proceeds with an activation barrier of 0.59 eV, which is significantly reduced with respect to the bulk water case (0.85 eV). Based on our results described above, however, we expect most of the  $\text{Ti}_{5c}$  sites to be occupied by  $\text{H}_2\text{O}$  molecules that are difficult to displace. We also observe that  $\text{H}_3\text{O}^+$  formed during this reaction immediately splits into an  $\text{H}_2\text{O}$  molecule and an  $\text{H}^+$  ion which is then trapped by the nearby bridging oxygen  $\text{O}_b$  atom (see Figure 7). Thus, our calculations suggest that some of the surface  $\text{O}_b$  sites may be protonated as the result of  $\text{HCO}_3^-$  formation. At low pH the  $\text{O}_b$  sites are expected to be already protonated, diminishing the probability of bicarbonate formation from  $\text{CO}_2$  adsorbed at  $\text{Ti}_{5c}$ , whereas high pH should facilitate the reaction. Also, at high pH the formation of bicarbonate will be also facilitated due to  $\text{CO}_2$  reaction with  $\text{OH}^-$  groups in solution, which is characterized by a much smaller free-energy barrier of 0.6 eV than between  $\text{CO}_2$  and  $\text{H}_2\text{O}$  (0.82 eV).<sup>21,55</sup>

Another reaction mechanism of  $\text{HCO}_3^-$  formation at the  $\text{TiO}_2$  surface has been recently hypothesized in experiments on  $\text{TiO}_2$  with bimetallic clusters.<sup>11</sup> This mechanism suggests the participation of  $\text{TiO}_2$  lattice oxygen atoms in oxidation of the protonated  $\text{CO}_2$  species at the surface. As discussed above, we indeed observe spontaneous protonation of  $\text{CO}_2$  in the  $\text{O}_b\text{-CO}_2$  configuration to yield  $^*\text{COOH}$  and thus explore the kinetics of the reaction  $\text{O}_b\text{-COOH} \rightarrow \text{HCO}_3^-$  (aq). We find, however, that this reaction pathway is characterized by a very high activation barrier of  $\sim 2$  eV.

In the following, we examine several plausible adsorption configurations of  $\text{HCO}_3^-$  on the  $\text{TiO}_2(110)$  surface. Among different possibilities, we determine two configurations that stay stable after 5 ps CPMD run. The first stable state is a  $\text{HCO}_3^-$ (1) molecule adsorbed in a monodentate way at  $\text{Ti}_{5c}$  (Figure 8a). The average distance between surface  $\text{Ti}_{5c}$  and O is found to be 2.201 Å and the free energy of desorption from this

configuration is estimated at 0.49 eV. Thus,  $\text{HCO}_3^-$  is predicted to be bound to the  $\text{Ti}_{5c}$  site considerably stronger than  $\text{CO}_2$ , but weaker than  $\text{H}_2\text{O}$ .

Figure 8 shows the second stable state of a  $\text{HCO}_3^-$ (2) molecule, which is bidentate adsorbed at two adjacent  $\text{Ti}_{5c}$  sites. The average  $\text{O-Ti}_{5c}$  distance is now 2.163 Å, which is slightly shorter than the  $\text{Ti-O}$  distance for the monodentate adsorbed configuration. The desorption process of bidentate configuration  $\text{HCO}_3^-$ (2) is considered in two steps: (a) breaking one of the two bonds and thereby switching the molecule to monodentate configuration  $\text{HCO}_3^-$ (1) and then (b) detachment of a monodentate  $\text{HCO}_3^-$ (1) molecule. The metadynamics simulation reveal that  $\text{HCO}_3^-$  is very strongly bound to two surface  $\text{Ti}_{5c}$  atoms and the free energy to break the first bond is 0.87 eV. Thus, the obtained results prompt that the formation of a very tightly bound bicarbonate monolayer on the rutile (110) surface could block all of the unsaturated Ti atoms and prevents all the described above reactions of  $\text{CO}_2$  reduction,<sup>15</sup> but only if strongly bound  $\text{H}_2\text{O}$  molecules had been displaced from  $\text{Ti}_{5c}$ .

## CONCLUSIONS

In summary, Car–Parrinello molecular dynamics in conjunction with metadynamics simulations were used to describe the competitive adsorption/desorption behavior of  $\text{CO}_2$  and  $\text{H}_2\text{O}$  over the rutile  $\text{TiO}_2(110)$  surface in water environment at room temperature. It was determined that in the aqueous phase  $\text{CO}_2$  preferentially adsorbs at the bridging oxygen ( $\text{O}_b$ ) atoms, while  $\text{Ti}_{5c}$  sites are saturated by strongly bound  $\text{H}_2\text{O}$  molecules. Our calculations reveal that  $\text{CO}_2$  is spontaneously protonated in the adsorbed state leading to the formation of activated  $^*\text{COOH}$  species at the surfaces, and the kinetics of  $\text{OH}^-$  detachment to form CO is greatly enhanced in the presence of a  $\text{Ti}^{3+}$  polaron. The calculations also suggest that if tightly bound  $\text{H}_2\text{O}$  molecules are displaced from  $\text{Ti}_{5c}$  sites, then the  $\text{TiO}_2(110)$  surface catalyzes the formation of surface-bound  $\text{HCO}_3^-$  species that will block the surface from further reactions.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.7b02777.

Detailed description of the atomistic model, Car–Parrinello molecular dynamics, and metadynamics simulations with all computational parameters (PDF)

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## Notes

The authors declare no competing financial interest.

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